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# Dark matter astrometry: accuracy of subhalo positions for the measurement of self-interaction cross-sections

David Harvey,<sup>1</sup>★ Richard Massey,<sup>1,2</sup> Thomas Kitching,<sup>1</sup> Andy Taylor,<sup>1</sup> Eric Jullo,<sup>3</sup> Jean-Paul Kneib,<sup>3</sup> Eric Tittley<sup>1</sup> and Philip J. Marshall<sup>4</sup>

<sup>1</sup>*SUPA, University of Edinburgh, Royal Observatory, Blackford Hill, Edinburgh EH9 3HJ, UK*

<sup>2</sup>*Institute for Computational Cosmology, Durham University, South Road, Durham DH1 3LE, UK*

<sup>3</sup>*Laboratoire d'Astrophysique de Marseille, CNRS-Université Aix-Marseille, 38 rue F. Joliot-Curie, F-13388 Marseille, France*

<sup>4</sup>*Department of Physics, University of Oxford, Keble Road, Oxford OX1 3RH, UK*

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## ABSTRACT

Direct evidence for the existence of dark matter and measurements of its interaction cross-section have been provided by the physical offset between dark matter and intracluster gas in merging systems like the Bullet Cluster. Although a smaller signal, this effect is more abundant in minor mergers where infalling substructure dark matter and gas are segregated. In such low-mass systems the gravitational lensing signal comes primarily from weak lensing. A fundamental step in determining such an offset in substructure is the ability to accurately measure the positions of dark matter subpeaks. Using simulated *Hubble Space Telescope* observations, we make a first assessment of the precision and accuracy with which we can measure infalling groups using weak gravitational lensing. We demonstrate that using an existing and well-used mass reconstruction algorithm can measure the positions of  $1.5 \times 10^{13} M_{\odot}$  substructures that have parent haloes 10 times more massive with a bias of less than 0.3 arcsec. In this regime, our analysis suggests the precision is sufficient to detect (at  $3\sigma$  statistical significance) the expected mean offset between dark matter and baryonic gas in infalling groups from a sample of  $\sim 50$  massive clusters.

**Key words:** gravitational lensing: weak – cosmological parameters – dark matter – galaxies: clusters general.

## 1 INTRODUCTION

Evidence for dark matter (DM) has been accumulating for 80 yr (e.g. Zwicky 1933; Rubin, Ford & Thonnard 1980; Battaner & Florido 2000; Clowe et al. 2006), yet its nature and properties remain poorly understood. How DM manifests itself in the Universe has become a key question in both particle and astrophysics, which has resulted in a variety of studies all attempting to shed some light on this dark mystery.

Despite evidence from accumulating from astronomical sources, the tightest constraints on the properties of DM are being led by terrestrial experiments. The Large Hadron Collider is attempting to create new particles in high-energy proton–proton collisions that could potentially be DM candidates (e.g. Baer & Tata 2009; Mitsou 2011). Furthermore, direct detection experiments are trying to observe the galactic DM wind caused by the orbit of the Solar system around the Galaxy, and the Earth around the Sun (e.g. Bertone,

Hooper & Silk 2005; Burgos et al. 2009; Bernabei et al. 2010; Angloher et al. 2012).

Astronomical techniques currently provide looser constraints on otherwise inaccessible parameters. DM annihilation signals could be observed through gamma-rays originating from the centre of the galaxy (e.g. Cholis et al. 2009; Hooper & Goodenough 2011), placing constraints on the annihilation cross-section. Alternatively the density profiles of galaxy clusters can constrain the self-interaction cross-section. Yoshida et al. (2000) found that a few collisions per particle per Hubble time can significantly affect the profile at the core of the cluster. Moreover, cosmological simulations using self-interacting DM have shown that a small but finite cross-section will have an effect on the core size and central density (Rocha et al. 2013). Peter et al. (2013) tried to compare these simulations with observations and found that constraints from such a technique will most probably be improved by measurements of central densities and not halo shapes. Both techniques provide a unique way to probe the properties of DM.

The trajectories of different mass components during major mergers like the Bullet Cluster (1E 0657–558) have recently provided

★ E-mail: drh@roe.ac.uk

important constraints on the DM–DM self-interaction cross-section  $\sigma$  (Clowe, Gonzalez & Markevitch 2004; Bradač et al. 2006; Clowe et al. 2006; Mahdavi et al. 2007). The baryonic components of a galaxy cluster can be seen via direct imaging: optical emission from galaxies, and X-ray bremsstrahlung emission from hot gas. DM cannot be observed directly, but its projected density can be reconstructed from the observable ‘gravitational lensing’ of the images of background sources behind the cluster (see reviews Bartelmann & Schneider 2001; Refregier 2003; Hoekstra & Jain 2008; Massey, Kitching & Richard 2010). In the model for the Bullet Cluster collision, the X-ray emitting gas in the intracluster medium was slowed during its first core passage, preventing it from travelling far from the point of impact. The DM in each cluster, interacting essentially only via gravity and mapped using its weak and strong gravitational lensing effects, is supposed to have passed through unaffected. In this picture, the temporary separation of the Bullet Cluster’s DM from its gas yields a constraint of  $\sigma/m < 1 \text{ cm}^2 \text{ g}^{-1}$  (Markevitch et al. 2004). Similar analyses yield  $\sigma/m < 4 \text{ cm}^2 \text{ g}^{-1}$  from cluster MACS J0025–1222 (Bradač et al. 2008) and  $\sigma/m < 3 \text{ cm}^2 \text{ g}^{-1}$  from cluster Abell 2744 (Merten et al. 2011).

As shown by Williams & Saha (2011) a further displacement can be measured between the DM and the stars of the cluster. Assuming that stars act as non-interacting tracer particles, any collision-induced separation between the DM density peaks (inferred from their lensing effects) and the cluster member galaxies (visible by their starlight) could provide evidence for there having been some self-interacting DM. The offset between DM and stars in cluster Abell 3827 intriguingly suggests a non-zero interaction cross-section, with lower limit  $\sigma/m > 4.5 \times 10^{-7} (t/10^{10} \text{ yr})^{-2} \text{ cm}^2 \text{ g}^{-1}$ , where  $t$  is the infall time of the subclump around the main halo (Williams & Saha 2011). This result is, however, sensitive to the interpretation of a very small number of proposed multiple (strongly lensed) images.

In this work we assess the precision and accuracy of weak gravitational lensing measurements of the *position* of mass peaks. This differs from the many studies that have assessed the precision and accuracy of measurement of the *mass* of mass peaks. We investigate whether it will be possible to detect small offsets in position on the sky between the baryonic and DM density peaks of cluster substructures (Massey, Kitching & Nagai 2011, hereafter MKN). We imagine detecting these infalling galaxy groups, and measuring their barycentres, from their X-ray (or optical) emission, and comparing with the positions of mass density peaks reconstructed by analysis of weakly lensed background objects in the vicinity. Such analyses have been carried out in individual interacting clusters using flexible exploratory mapping techniques by (e.g. Clowe et al. 2004; Markevitch et al. 2004); here we consider measuring offsets – ‘bulleticities’ – in many different clusters, and combining the results in a statistical measurement of the interaction cross-section (MKN). In particular, we are interested in using analytically simulated data to answer the following questions.

- (i) To what precision can we measure the offset in a single infalling substructure?
- (ii) Can we identify a point estimator whose simple combination over a sample will provide a measurement of subhalo position with minimal bias?
- (iii) What are the dominant sources of residual bias in this estimate?
- (iv) How large a sample of observed clusters are we likely to need to be able to detect an offset between DM and baryonic as predicted by MKN?

- (v) What further investigation might be needed to prove the utility of this technique for probing DM interaction cross-sections?

This paper is organized as follows. In Section 2 we outline the theory behind weak gravitational lensing and our goals in its application. In Section 3 we present an end-to-end simulation pipeline in which we start with a known mass distribution, simulate *Hubble Space Telescope* (HST) lensing data, then use LENSTOOL to reconstruct the mass distribution. In Section 4 we describe our results. In Sections 6 and 7 we conclude and outline future work.

## 2 THEORY

### 2.1 Weak gravitational lensing

Gravitational lensing is the deflection of light rays by the distortion of space–time around any massive object. This phenomenon can be used to map the distribution of mass, including otherwise invisible DM.

A 3D density distribution  $\rho$  with gravitational potential  $\Phi$  can be projected on to the plane of the sky to obtain a 2D deflection potential:

$$\Psi \equiv \frac{D_{\text{OL}} D_{\text{LS}}}{D_{\text{OS}}} \frac{2}{c^2} \int \Phi \, dz, \quad (1)$$

where the angular diameter distances between the observer, lens and source encode the geometry of a convergent lens. The image of a background galaxy passing through this potential is magnified by a convergence

$$\kappa = \frac{1}{2} \left( \frac{\partial^2 \Psi}{\partial x^2} + \frac{\partial^2 \Psi}{\partial y^2} \right) \quad (2)$$

and distorted by a shear

$$\gamma_1 = \frac{1}{2} \left( \frac{\partial^2 \Psi}{\partial x^2} - \frac{\partial^2 \Psi}{\partial y^2} \right), \quad (3)$$

$$\gamma_2 = \frac{\partial^2 \Psi}{\partial x \partial y}, \quad (4)$$

where  $\gamma_1$  ( $\gamma_2$ ) refers to elongation along (at  $45^\circ$  to) an  $x$ -axis defined arbitrarily in the plane of the sky (Bartelmann & Schneider 2001; Refregier 2003).

Around a foreground galaxy cluster, background galaxies appear aligned in distinctive circular patterns, with tangential and curl components:

$$\gamma_t = -[\gamma_1 \cos(2\phi) + \gamma_2 \sin(2\phi)], \quad (5)$$

$$\gamma_\times = -\gamma_1 \sin(2\phi) + \gamma_2 \cos(2\phi), \quad (6)$$

where  $\phi$  is the angle of the galaxy position with respect to the Cartesian axis.

If galaxies were intrinsically circular and of fixed size, the applied shear would simply change their apparent ellipticity. In practice it is necessary to average the observed shapes of  $\sim 100$  galaxies to remove the influence of their complex morphologies; and because of a degeneracy between shear and magnification, only the ‘reduced shear’

$$g \equiv \frac{\gamma}{1 - \kappa} \quad (7)$$

is observable.

## 2.2 Bulleticity

Attempts to constrain the self-interaction cross-section of DM from major galaxy cluster collisions face two obstacles. First, measuring a separation between the DM and baryonic gas requires a merger between clusters of similar masses to be seen at just the right time since first core passage, and these are rare events (Shan, Qin & Zhao 2010). Second, uncertainties in the impact velocity, impact parameter and orientation with respect to the line of sight severely limit constraints from individual clusters (Randall et al. 2009).

Fortunately, hydrodynamical simulations of structure formation (Nagai et al. 2007a,b; Powell, Kay & Babul 2009) predict a mean offset between DM and gas during the infall of *all* subhaloes into massive clusters (MKN). Although weak gravitational lensing cannot precisely resolve the positions and masses of individual pieces of small substructure, a statistical ‘bulleticity’ signal can be obtained by averaging the measurements from many clusters. The bulleticity vector  $\mathbf{b}$  is the offset between substructure’s total mass (where DM dominates) and baryonic components in the plane of the sky:

$$\mathbf{b} \equiv b_r \hat{\mathbf{e}}_r + b_t \hat{\mathbf{e}}_t, \quad (8)$$

where  $\hat{\mathbf{e}}_r$  and  $\hat{\mathbf{e}}_t$  are unit vectors in the radial and tangential directions with respect to the cluster centre. Hydrodynamical simulations show that, despite complex and interacting processes, the net effect of cluster gas physics is a force on the substructure gas similar to a simple buoyancy that produces an offset  $\langle b_r \rangle > 0$ . This is the key signal in which we are interested. The simulations also show that, with a sufficiently large sample and no preferred in-fall handedness,  $\langle b_t \rangle \equiv 0$ . Checking that measurements of this are consistent with zero will be a useful test for residual systematics.

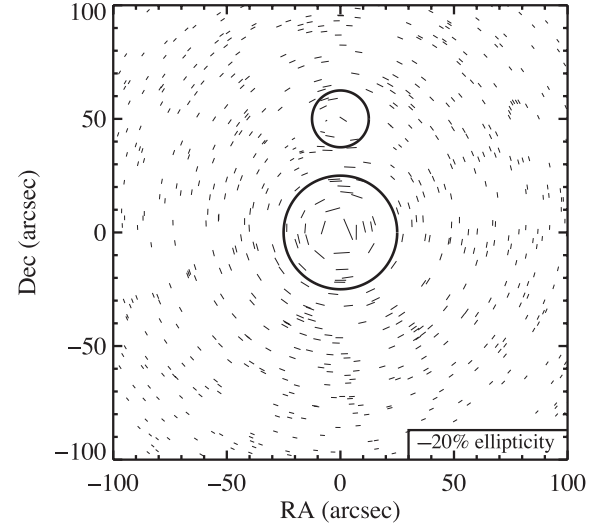
MKN showed that for a  $\Lambda$  cold dark matter ( $\Lambda$ CDM) paradigm with collisionless DM, the expected radial offset between baryonic and dark components of substructure is  $\sim 10, 3.5$  and  $2$  arcsec at a redshift  $z = 0.1, 0.3, 0.6$ , respectively at a radial distance of  $0.15r_{500}$ , which increased towards the centre of the cluster. Therefore the measurement of an offset relies upon an ability to measure the position of substructure components with minimal bias near to the core. Statistical errors will be gradually beaten down by averaging many measured offsets. However any systematic bias in the centroid of either component will propagate into constraints on the interaction cross-section.

The spatial resolution of the X-ray space telescope *Chandra* is subarcsecond, whereas any weak lensing map will be limited by the finite density of resolved galaxies to  $\sim 10$  arcsec with the deepest, highest resolution data (e.g. Massey et al. 2007a). Although the accuracy in which one can define the X-ray peak will depend on the distribution of gas the DM peak will be affected by a similar problem. Therefore it is possible to assume that the error in the X-ray peak position is subdominant to that of DM, and study in detail the reliability of weak lensing centroid estimates only. We will also quantify the precision of weak lensing centroiding to estimate the sample size required to detect bulleticity.

## 3 METHOD

### 3.1 Simulated shear fields

In order to examine the exact behaviour of weak lensing as a positional estimate of DM we need initial experiments in carefully controlled environments. This includes having mass distributions with well-defined correct answers rather than cosmological simulations. We therefore create simulated shear fields containing DM haloes



**Figure 1.** The reduced shear signal of a simulated cluster with a NFW profile. The main halo has a  $M_{200}$  of  $8 \times 10^{14} M_\odot$  and is positioned at (0,0) and the subhalo,  $8 \times 10^{13} M_\odot$ , is positioned at (0,49). The field of view represents that of a *HST/ACS* with a typical density of galaxies of  $80 \text{ arcmin}^{-2}$ . The circles are a guide for where they are placed and have no physical significance.

of known position, mass and ellipticity. For an analytic model, we adopt the Navarro, Frenk & White (1996, hereafter NFW) density profile for a galaxy cluster at a conservative redshift of 0.6,

$$\rho(r) \propto \frac{1}{\frac{r}{r_s} \left(1 + \frac{r}{r_s}\right)^2}, \quad (9)$$

where the scale radius  $r_s$  can be expressed in terms of the concentration parameter  $c = r_{\text{vir}}/r_s$  (and  $r_{\text{vir}}$  is the virial radius). For typical clusters, an empirical relation (Macciò, Dutton & van den Bosch 2008) suggests

$$\log \langle c \rangle = 0.830 - 0.098 \log (M_{\text{vir}} / [10^{12} h^{-1} M_\odot]), \quad (10)$$

where  $M_{\text{vir}}$  is the virial mass.

Using analytical equations for multiple NFW haloes (Lasky & Fluke 2009; White & Brainerd 2000), we construct a cluster system with infalling galaxy group(s). As shown in Fig. 1, the baseline configuration includes a main halo in the centre of the field of view plus a subhalo 49 arcsec to the north. Substructure typically contains  $\sim 10$  per cent of the mass of a system (e.g. Cohn 2012), we fix the mass of the subhalo to be always 10 per cent that of the parent halo (with a concentration parameter given by equation 10).

The shears due to multiple mass components in the same cluster simply add, such that  $\gamma_{\text{total}} = \gamma_{\text{main halo}} + \gamma_{\text{subhalo}_1} + \gamma_{\text{subhalo}_2} + \dots + \gamma_{\text{subhalo}_n}$ .

In order to get high signal-to-noise ratio imaging of cluster cores we manufacture simulated *HST* weak lensing measurements of the cluster system. In a  $200 \times 200 \text{ arcsec}^2$  field of view, shear measurements are simulated from 80 (randomly placed) galaxies per square arcminute, as could be obtained from a full one-orbit exposure in the *F814W* band with the Advanced Camera for Surveys (ACS). In a typical Hubble archive the redshift of a cluster is varies between 0 and 1. We choose a conservative redshift of 0.6, at which the bulleticity signal is small but potentially detectable. For each cluster configuration and mass, we generate 100 noise realizations.

### 3.1.1 Intrinsic galaxy morphologies

The basic challenge with weak lensing measurements is that galaxies are not inherently circular – indeed, the ellipticity of a typical galaxy is an order of magnitude larger than the shear. Chance alignments of galaxies can mimic a coherent gravitational lensing signal and introduce noise into the mass reconstruction. In individual clusters, such noise can cause centroid shifts of  $\sim 10$  arcsec, but this should average away as long as the noise has no preferred direction over a sample of clusters. Leauthaud et al. (2007) state that galaxies typically have an intrinsic ellipticity distribution with a mean of zero and a width,  $e_{\text{int}}$ , of 0.3. This can be expressed as a complex number by  $e_{\text{int}} = |e_{\text{int}}| \exp(2i\theta)$ , where  $\theta$  is the angle of the galaxy. One can then transform the galaxy from the source plane to the image plane using the complex reduced shear,  $g = |g| \exp(2i\phi)$ , where  $\phi$  is orientation of the galaxy due to the lensing effect, via,

$$e^{(I)} = \frac{e_{\text{int}} + 2g + g^2 e_{\text{int}}^*}{1 + |g|^2 + 2\Re(g e_{\text{int}}^*)}, \quad (11)$$

where the star represents the complex conjugate.

### 3.1.2 Elliptical mass distributions

Galaxy clusters are often not spherically symmetric (Jetzer et al. 2001). Misidentifying the shape of a halo can introduce spurious detections of substructure along the major axis, or shift the apparent position of real substructure. It is therefore important to check whether elliptical haloes affect the centroid estimate of both the cluster and the subhalo. We have run simulations with both a spherical and an elliptical main halo. In elliptical cases, the ellipticity of the main potential was fixed at 0.2 (where ellipticity  $= [a^2 - b^2]/[a^2 + b^2]$ ). To span a range of possible scenarios, the major axis is aligned at  $0^\circ$ ,  $45^\circ$  or  $90^\circ$  from the positive  $x$ -axis (in the latter case, this points towards the subhalo).

It was considered that force fitting a circularly symmetric fit to an elliptical main halo could potentially bias the position however the signal-to-noise ratio of the subhalo would mean that constraining the ellipticity would not be possible. Moreover DM would interact only gravitationally, and therefore we expect any infalling halo to retain its radial symmetry, thus in all cases the subhalo is kept circular.

### 3.1.3 Imperfect shape measurement

Achieving sub-per cent accuracy in the measurement of galaxies' apparent shapes is an ongoing challenge (Heymans et al. 2006; Massey et al 2007b; Bridle et al. 2010; Kitching et al. 2012a,b). Even with a space-based telescope, the point spread function (PSF) varies across the field of view and can change over time (Rhodes et al. 2007). If the PSF is not accurately modelled or the image effectively deconvolved, it can be spuriously imprinted upon the shear measurements. Image noise and pixelation further impede the measurement of small, faint galaxy shapes.

We do not consider multiplicative shear measurement biases here, since they bias only the recovered mass estimates, and not the positions. We do, however, consider additive shear measurement biases, which will affect the inferred mass clump positions. The PSF normally has a preferred direction with respect to the telescope, but the location of substructure and the angle of orientation at which the cluster is imaged will vary from cluster to cluster. In each realization of a simulated catalogue, we add a constant spurious signal  $c_i$  to each component of shear, drawn from a Gaussian distribution with mean

0 and standard deviation 0.01 (which is split into shear components so that it is in a random direction). Finally, we model the pixelation noise by adding an additional stochastic component to each shear measurement, drawn from a Gaussian distribution with mean 0 and width 0.01 again split into components, although this is effectively degenerate with (and subdominant to) the intrinsic ellipticity. Thus the observed shears become

$$e_i^{\text{tot}} = e_i^{(I)} + c_i + \sigma_i^N. \quad (12)$$

Although many algorithms linearize the lensing potential, which does not hold in regimes of  $g > 0.1$ , the bias introduced by this assumption can be considered a multiplicative factor to the shear and in fact would not have an effect to the position of the subhalo (Massey & Goldberg 2008; Melchior & Viola 2012).

### 3.1.4 Galaxy redshift distribution

A statistical measurement of bulleticity will require a large sample of galaxy clusters, and multicolour imaging may not be available for them all. The distortion experienced by each galaxy image depends upon the lensing geometry as described by equation (1). With only monochromatic imaging, it can even be impossible to tell whether galaxies are behind a cluster (and therefore lensed) or in front of it (and therefore undistorted). Allowing foreground galaxies in the galaxy catalogue will dilute the inferred shear signal. We introduce a source galaxy redshift distribution:

$$p(z) \propto z^2 \exp \left[ - \left( \frac{z}{z^*} \right)^{1.5} \right], \quad (13)$$

where  $z^* = z_{\text{med}}/1.1412$  and  $z_{\text{med}} = 1.0$  (Taylor et al. 2007). We apply this uniformly across the field of view. We assume we know exactly the redshift of the cluster (0.6) and for the purposes of measuring the position we do not concern ourselves with  $\Sigma_{\text{crit}}$  and the total mass.

A further problem with having only monochromatic imaging is that it will be impossible to distinguish between background sources and cluster members. Although the effect of this would be further dilution of the signal, since the members will be correlated with the density profile of the cluster the dilution will also be correlated. It is therefore possible for the position of the halo to be biased if these galaxies are included in the reconstruction. Therefore, a simple distribution of member galaxies is placed over the cluster such that they follow the NFW profile. The number of member galaxies is then increased and to study the effect the member galaxies may have.

### 3.1.5 Multiple substructures

In the paradigm of hierarchical structure formation, clusters grow through multiple mergers, so multiple subhaloes may be physically close to a cluster at a given time. The presence of multiple subhaloes will complicate the shear field and thus make it harder to estimate the positions of each. It is therefore important to be confident that if subhaloes are close together in real space their signals do not cause a bias in any direction. A set of realizations were run with a second subhalo was introduced into the field. The first subhalo remained at (0, 49) arcsec from the main halo; to span a range of possible configurations, the second halo was placed at three different positions (49, 49), (−49, 0) and (0, −49) arcsec from the main halo.



### 3.1.6 Potential line of sight contamination

Independent large-scale structure at different redshifts may happen to lie along the line of sight to the cluster, and be misinterpreted as substructure (Hoekstra 2001; Spinelli et al. 2012). As unassociated galaxy groups will not be falling into the cluster, they will not exhibit any systematic offset between DM and gas, and their inclusion will spuriously dilute the measured bulleticity. Since substructure will be initially identified via X-ray imaging, we can estimate the number of coincidentally aligned structures by considering the density of X-ray luminous groups in unpainted observations.

In the 1.64 deg<sup>2</sup> Cosmic Evolution Survey (COSMOS; Scoville et al. 2007), Finoguenov et al. (2007) found 206 X-ray groups with masses  $10^{13} - 10^{14} M_{\odot} h_{72}^{-1}$ , which matches the mass range considered in this work. Leauthaud et al. (2010) tried to detect all of these groups via weak lensing from one-orbit *HST* imaging. About a quarter of the groups are detected at greater than our  $2\sigma$  detection threshold. Scaling this down to the  $200 \times 200$  arcsec<sup>2</sup> ACS field of view, we expect a contaminant of around one spurious peak for every 20 clusters. This  $\sim 5$  per cent dilution should be considered in a second larger survey [e.g. Hyper Suprime-Cam (HSC), Dark Energy Survey (DES), *Euclid*], and could be reduced if (even coarse) photometric redshifts were available for some galaxies.

### 3.1.7 Substructure as a function of distance from the cluster centre and mass ratio

The basis simulation is set up such that the distance the subhalo is from the centre of the cluster and the mass fraction between the main halo and the subhalo is held constant. Although the values used for the simulations are that of a typical cluster these will not be constant in the case of real data and therefore the radial distance and the mass fraction are both independently varied.

## 3.2 Mass reconstruction

Many algorithms have been developed to reconstruct the mass, concentration and position of massive ( $> 10^{14} M_{\odot}$ ) haloes from observations of weak (and strong) gravitational lensing (Bradač et al. 2005; Cacciato et al. 2006; Diego et al. 2007; Merten et al. 2009; Zitrin et al. 2012). However, testing of these has generally focused on the mass and concentration parameters, positional accuracy has not yet been pushed to the low  $\sim 10^{13} M_{\odot}$  mass regime.

To determine the viability of bulleticity measurements, weak lensing reconstructions must be tested in scenarios that reflect the environments in which it will be used. The main requirement for bulleticity is an accurate estimate of the subhalo and main halo positions with minimal bias.

LENSTOOL (Jullo et al. 2007) is open source software that calculates analytical models of the lensing signal for specific cluster density profiles and then compares them against the observed data. In a Bayesian framework, LENSTOOL samples the posterior using a Markov chain Monte Carlo algorithm. It continually probes the entire parameter space providing an estimate of entire posterior surface. The posterior for a given prior,  $\Pi$ , and likelihood,  $L$ , is

$$P = L\Pi, \quad (14)$$

where the likelihood is Gaussian,

$$L = \frac{1}{(2\pi\sigma^2)^{N/2}} e^{-\chi^2/2}. \quad (15)$$

The  $\chi^2$  statistic in LENSTOOL is calculated by converting the observed galaxy ellipticity to the source plane using the proposed model parameters. The resulting ellipticity should represent the intrinsic shape of the galaxy, which when summed over the entire field will have a mean of zero with some known variance. A chi-squared close to one shows that the model parameters to convert the image to the source plane were a good fit. Thus the chi-squared for given set of parameters, calculated in the source plane, is simply

$$\chi^2 = \sum_{i=1}^N \sum_{j=1}^2 \frac{(e_{i,j}^{(s)})^2}{\sigma_{i,j}^2}, \quad (16)$$

where the total error in the ellipticity,  $\sigma$ , is the sum of the intrinsic ellipticity and the shape measurement error added in quadrature, i.e.  $\sigma = \sqrt{\sigma_{\text{int}}^2 + \sigma_{\text{shape}}^2}$ .

The fact that the chi-squared is calculated in the source plane means the reconstruction can be affected by the way LENSTOOL converts from the source plane to the image plane. The input parameters for LENSTOOL are the semimajor axis,  $a$ , the semiminor axis,  $b$ , and the angle of the galaxy with respects to the image  $x$ -axis. These ellipse descriptors not only define the ellipticity of the galaxy but also the size. It has been shown in Schmidt et al. (2012) that measuring the sizes of galaxies is difficult and also ambiguous in how one defines it therefore we would like to avoid using this parameter.

We therefore decide to use *option 7* in LENSTOOL,<sup>1</sup> which transforms the  $a$ ,  $b$  and angle of the galaxy into the complex ellipticity  $e^{(l)}$  in the image plane using

$$e^{(l)} = \frac{a^2 - b^2}{a^2 + b^2} \exp(2i\theta), \quad (20)$$

removing any information on the size of galaxy. The ellipticity  $e^{(l)}$  is then transformed into the source plane via in the inverse of equation (11). From this the chi-squared calculated in equation (16) is made. Since LENSTOOL calculates the source ellipticity in this way we assign some nominal value to the size of the galaxy in the simulations.

Weak lensing mass reconstructions inevitably have limited resolution, because shear is a non-local effect (see equations 2–4) and because the shear field is sampled only at the positions of a finite number of background galaxies.

Fortunately, all that is required to get a robust measurement of bulleticity is unbiased centroid measurements. Where

<sup>1</sup> We also tested *option 6* which takes  $a$ ,  $b$  and the angle and transforms them directly to the source plane via

$$Q^{(s)} = AQA, \quad (17)$$

where  $Q$  is the quadrupole moment matrix,

$$Q_{ij} = \frac{\int d^2\theta q_l [I(\vec{\theta})](\theta_i - \Theta_i)(\theta_j - \Theta_j)}{\int d^2\theta q_l [I(\vec{\theta})]}, \quad i, j \in \{1, 2\}, \quad (18)$$

with  $q_l$  acting as a weight function that causes the integral to converge,  $\theta$  are the coordinates on the plane of the sky and  $\vec{\theta}$  is the centre of light. From this the  $e^{(s)}$  can be found via

$$e^{(s)} = \frac{Q_{11}^{(s)} - Q_{22}^{(s)} + 2iQ_{12}^{(s)}}{Q_{11}^{(s)} + Q_{22}^{(s)}}, \quad (19)$$

and therefore using equation (16) find the chi-squared. This requires full knowledge of the size of the galaxy in order to obtain correct  $a$ ,  $b$  and angle parameters. Failure to do so will cause a bias in the parameter estimation hence why we used *option 7*. Incidentally we found no increase in error by using *option 7*.

**Table 1.** Input values and priors used during the reconstruction on the main halo and subhalo in the simulations. The values in the square brackets refer to the range, and the dots refer to the different mass scales of the simulations.

	Input value	Prior	Type
Main position (arcsec)	(0, 0)	30 on (0, 0)	Gaussian
Sub 1 position (arcsec)	(0, 49)	25 arcsec on (0, 49)	Flat circle
Sub 2 position (arcsec)	(49, 49), (49, 0), (0, -49)	25 arcsec radius	Flat circle
Main halo mass ( $M_{\odot}$ )	$(1, 1.5 \dots 7.5, 8) \times 10^{14}$	$[0.5, 49] \times 10^{14}$	Flat
Subhalo mass ( $M_{\odot}$ )	$(1, 1.5 \dots 7.5, 8) \times 10^{13}$	$[0.5, 49] \times 10^{13}$	Flat
Main halo concentration	Mass:conc rel	[1, 10]	Flat
Subhalo concentration	Mass:conc rel	[1, 10]	Flat
Mass priors	$M_{\text{main}} > M_{\text{sub}}$	$M_{\text{main}} > M_{\text{sub}}$	Statement

available, strong lensing dramatically tightens the resolution of mass maps – but to rely on strong lensing would unacceptably reduce the number of clusters that we could use.

Since bulleticity measurements will always require overlapping X-ray observations, we will use them to inject information into the reconstruction as a Bayesian prior. By assuming that each X-ray peak has an associated group of galaxies, and that the maximum signal for bulleticity is  $\sim 10$  arcsec at a redshift of 0.1 (MKN), it is only necessary to consider mass peaks within a small area around the substructure and main cluster only, ignoring the rest of the field of view. So if the X-ray suggests a two-body configuration then this is the model that is used.

Table 1 shows the positions and masses of the clumps simulated with the associated priors used.

### 3.2.1 Estimation of subhalo position

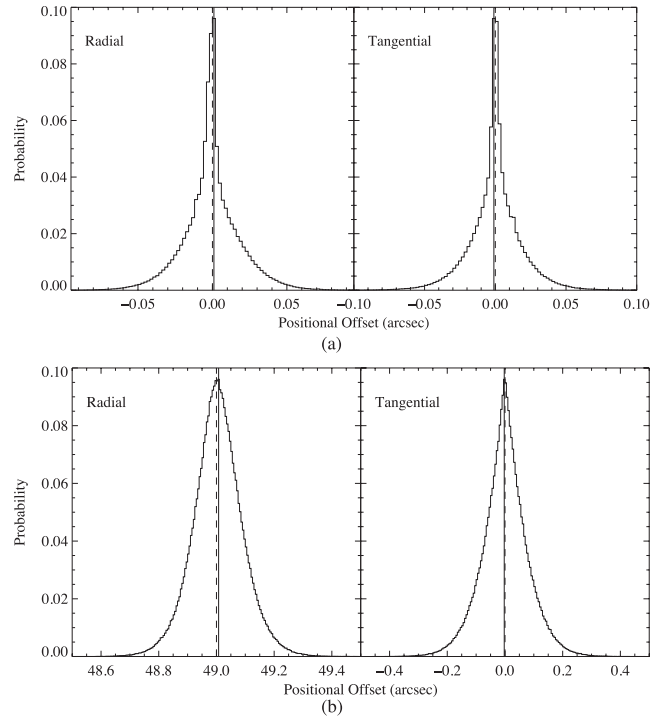
In order to understand LENSTOOL and its behaviour in the weak lensing limit for a two-halo system, we tested it on noise free simulations where the galaxies were inherently circular and the only affect was gravitational shear. Since LENSTOOL is a maximum likelihood algorithm in the case of zero noise the chi-squared calculation becomes undefined, so therefore we set the variance of ellipticity in LENSTOOL to a very small value (0.01).

Fig. 2 shows the full posteriors for the positions of the main halo and subhalo. The positions from the sampler have been binned with the maximum likelihoods shown as solid lines and the true values as dotted. The top panels show radial and tangential position of the main halo and the lower panels show the radial and tangential positions of the subhalo. It is clear that in the situation where there is no noise and the exact profile is given to LENSTOOL the maximum likelihood is centred on the true position with extremely small variance.

## 4 RESULTS

The expected offset between dark and baryonic components is  $\sim 2$  arcsec ( $\sim 3.5$  arcsec) at a redshift of 0.6 (0.3) at a radial of distance of  $0.15r_{500}$  and therefore any bias needs to be subdominant in comparison. The redshift distribution of clusters in the COSMOS field (Finoguenov et al. 2007) suggests that we expect a similar number of clusters at a redshift of 0.3–0.6, therefore, the measurement of an offset can tolerate  $\sim 0.5$  arcsec bias in the reconstruction in order to measure a bulleticity signal to  $\sim 3\sigma$  significance detection.

In an attempt to understand the behaviour of LENSTOOL, initial simple simulations were run and then an increasing number of contaminants and complexities were introduced. Unless stated otherwise, each panel in each figure shows  $\delta$ : the maximum likelihood



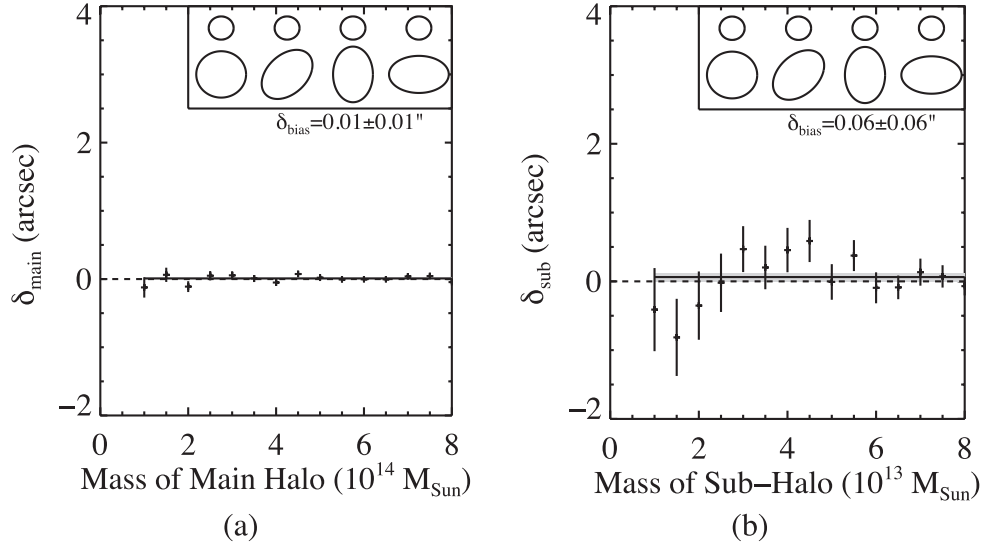
**Figure 2.** The likelihood surface for the main (top) and sub (bottom) halo positions in the case of zero noise (gravitational shear only). The binned histograms show the true posteriors and their maximum likelihoods as the solid line. The dotted line is the true position. The left-hand panels are the position in the radial direction and the right-hand panels are tangential direction. In the case of no noise the likelihood surface derived from LENSTOOL exhibits no bias around the maximum likelihood. (a) Main halo positional estimates. (b) Subhalo positional estimates.

radial position minus the true position for a given mass and simulation configuration, weighted averaged over 100 realizations for a given mass scale and then averaged over each configuration shown in the cartoon inset of the plot, i.e.

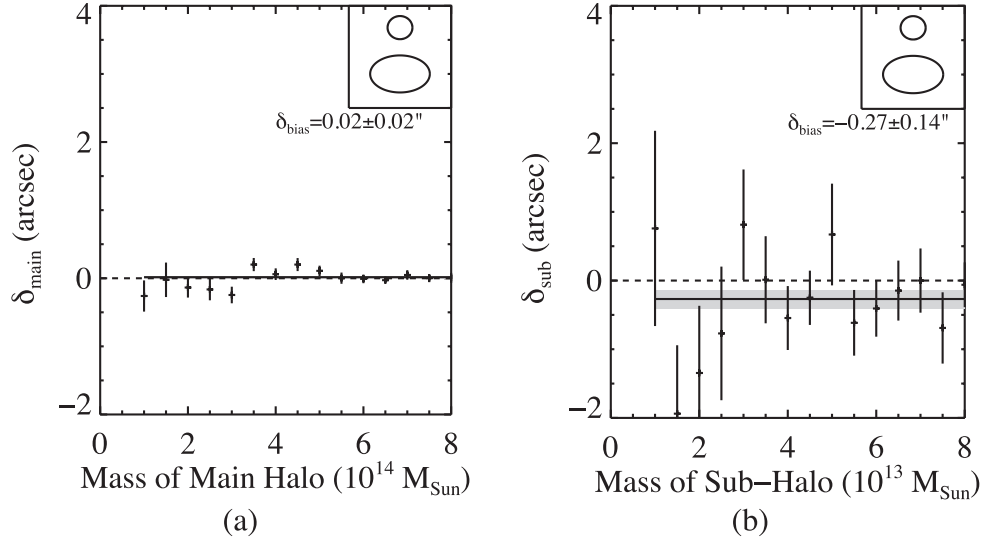
$$\delta_{\text{halo}}(m) = \langle (r_{\text{Meas}}(m) - r_{\text{True}})_{100} \rangle_{\text{config}}, \quad (21)$$

and the error for the given configuration is just given by the error in the mean. Furthermore, the main halo is always 10 times more massive than the subhalo (and appropriate concentrations given by equation 10).

We carried out four initial tests in the simplest two body case: one with a circular main halo and three with elliptical main haloes at different angles. Each test contained a simple background galaxy Gaussian intrinsic distribution and was run 100 times with different



**Figure 3.** Intrinsic ellipticities only: (a) and (b) show the positional estimates of the main halo and subhalo, respectively. In this initial test the background galaxies only contained a Gaussian intrinsic ellipticity distribution. The mass of the respective haloes is shown, in all cases the subhalo was 10 times less massive than the main halo [so results at  $8 \times 10^{14} M_{\odot}$  main halo in (a) are from the same simulation as those shown at  $8 \times 10^{13} M_{\odot}$  in (b)]. A variety of configurations were tested with the cartoon inset showing the set-up in each case. For each configuration, 100 noise realizations were run at the mass scale, the position of each halo estimated and then averaged over the all configurations (so each point reflects 400 averaged simulations). (a) Main halo positional estimates. (b) Subhalo positional estimates.



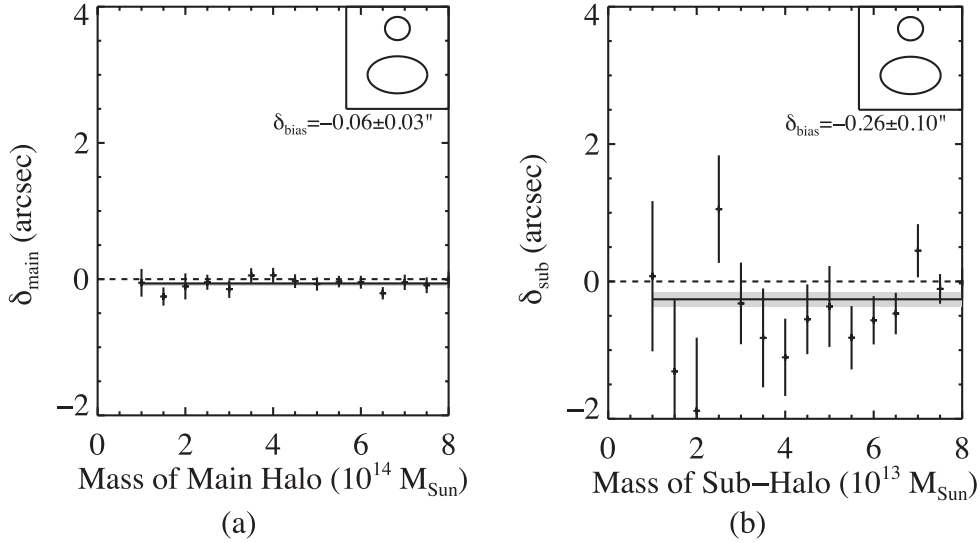
**Figure 4.** Intrinsic ellipticities and shape measurement bias: (a) and (b) are the positional estimates of the main halo and subhalo, respectively. In each case the mass is given and the main halo is 10 times more massive than the subhalo. The main halo is elliptical and the background galaxies have shape measurement bias and intrinsic ellipticities. In this case 100 realizations were run and the average position at each mass scale calculated. (a) Main halo positional estimates. (b) Subhalo positional estimates.

noise realizations. We then fitted two lines of best fit to the data to determine any significant bias in the positional estimates. One was a constant offset and the other a mass dependant one. For the subhalo we found that the reduced chi-square for a mass-dependant line was 1.34, whereas for a constant offset we found a chi-square of 1.25. Thus we found no significant evidence for a mass-dependant bias and therefore fitted a constant offset.

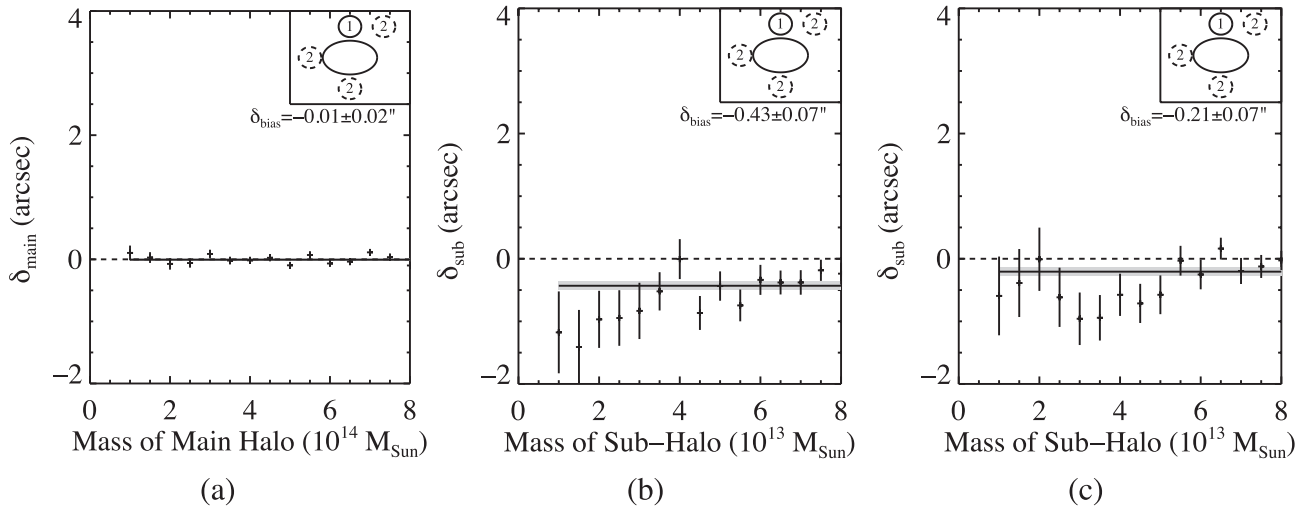
Each configuration showed in the cartoon inset of Fig. 3 exhibited no bias and therefore in order to better constrain the error on positional estimates we compiled the results into Fig. 3 giving the combined results from the initial tests.

It was found that LENSTOOL was robust to a basic level of noise so we introduced further sources of contaminants. Fig. 4 shows  $\delta$  for the main halo and subhalo, respectively, when shape measurement bias is introduced. Fig. 4(a) seems to show that the shape measurement bias has no affect on the positional estimate of the main halo, however, the subhalo in Fig. 4(b) seems to be slightly biased in the negative radial direction (towards the main halo). The cause of this will be the preferred direction of each galaxy. The level of the bias is of order 0.01, which is a similar level to the expected signal from a DM subhalo. Because each galaxy has a preferred direction it will mean that the preferred fit of the halo will not be the correct one,





**Figure 5.** Intrinsic ellipticities, shape measurement bias and source galaxy redshift distribution: (a) and (b) are the positional estimates of the main halo and subhalo, respectively. In each case the mass is given and the main halo is 10 times more massive than the subhalo. The main halo is elliptical and the background galaxies have shape measurement bias, a distribution in their redshift and intrinsic ellipticities. In this case 100 realizations were run and the average position at each mass scale calculated. (a) Main halo positional estimates. (b) Subhalo positional estimates.



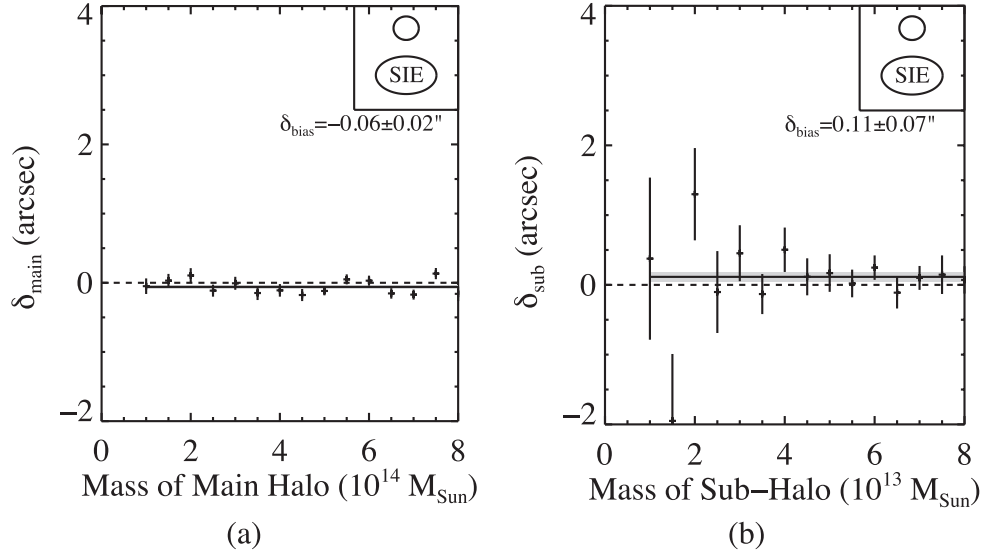
**Figure 6.** Dual subhalo simulation, with intrinsic ellipticities, shear measurement bias and source galaxy redshift distribution: (a) shows the positional estimates of the main halo, (b) shows the estimates of subhalo 1 and (c) gives the estimates of subhalo 2. The masses of the haloes are given. In each case the subhaloes are 10 times smaller in mass (so they are equal size) than the main halo. The background galaxies have intrinsic ellipticities, shape measurement bias and a redshift distribution. The plots show three different configurations (given by the dashed circles). In each case subhalo 1 is kept in the same place as shown in the cartoon inset, and for each of the three scenarios subhalo 2 is positioned as shown. In each scenario 100 noise realizations are run and the positions averaged over all configurations and noise realizations. (a) Main halo. (b) Subhalo 1. (c) Subhalo 2.

causing a bias in the position. This bias of  $0.27 \pm 0.14$  arcsec is well within the tolerated level.

Gravitational lensing is a geometrical effect and hence the signal is dependent on the distance the galaxy is from the halo. *LENSTOOL* requires knowledge of the source galaxy redshift, we therefore introduce a redshift distribution into to source galaxies and test the approximation that their redshifts are 1. Fig. 5 shows the results when such a distribution is introduced. Fig. 5(b) has no significant evidence for an increase in bias due a source galaxy redshift approximation from that of Fig. 4(b).

Using the same signal contaminants as Fig. 5 (intrinsic ellipticities, source redshift and shape measurement bias), we introduce a second subhalo, complicating the geometrical set-up of the simula-

tions. We ran three different scenarios; in each case the main halo was at the centre of the field and subhalo 1 was kept at the position previously simulated. The new, second subhalo was placed at three different locations as shown in the cartoon inset. For each subhalo 2 position 100 noise realizations were run. In all cases the main halo was 10 times more massive than the subhaloes, and the subhaloes were equal size. We found that there was no preferred bias dependant on the position of the second subhalo and so averaged these simulations together in order to better constrain the bias and error bars. Fig. 6 shows that the bias introduced by the source galaxy distribution is evident in the two-body system. The more complicated geometrical set-up seems to have no effect on the overall bias.



**Figure 7.** Simulated SIE haloes, with intrinsic ellipticities, source galaxy redshift distribution and shape measurement bias: (a) and (b) show the main halo and subhalo positional estimates, respectively. In this scenario a SIE profile is simulated and NFW fitted imitating profile misidentification in real data. The main halo is always 10 times larger than the subhalo, and the source galaxies have shape measurement bias, intrinsic ellipticities and a redshift distribution. 100 noise realizations were run and the average position estimated in each case. (a) Main halo positional estimates. (b) Subhalo positional estimates.

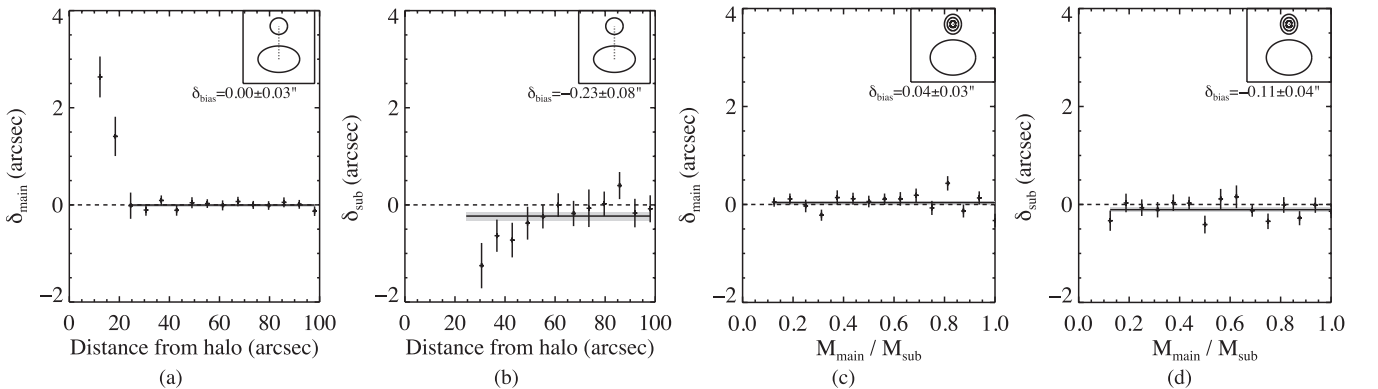
Fig. 7 is a test into the model dependency of the reconstruction method. Although NFW profiles have been extensively studied with both simulated and empirical data, the inclusion of baryons has shown that profiles depart from the assumed NFW (Duffy et al. 2010). Therefore in order to understand what the effect of this is, Fig. 7 shows the positional accuracy of the main halo and subhalo in the case where a singular isothermal ellipsoid (SIE) is simulated and using LENSTOOL a NFW is fitted. Fig. 7(b) shows that the subhalo has an insignificant positive bias, seemingly in contradiction to previous results.

This unbiased nature is due to the fact that the central core of an SIS is extremely peaked. An NFW has a much flatter profile in the core and therefore may introduce more uncertainty in the peak position. One effect of introducing an SIS is that the scatter seems to be much larger and at smaller masses the positional estimates become unreliable.

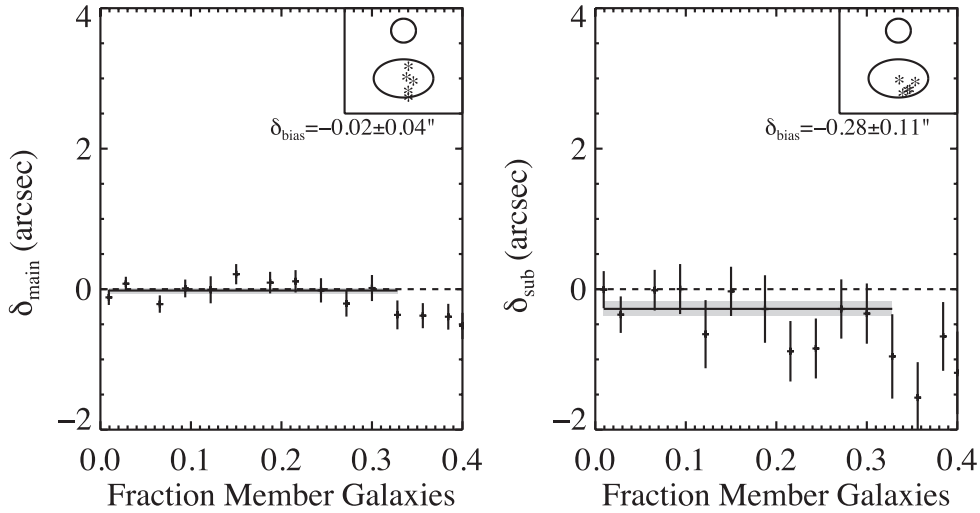
#### 4.1 Accuracy as a function of distance from the cluster and mass fraction

Fig. 8 shows the results if the halo masses are kept constant ( $8 \times 10^{14} M_{\odot}$ ) at a redshift of 0.6 with an ellipticity of 0.2, positional angle of  $180^\circ$ , shape measurement bias and a background galaxy redshift distribution, and (first two panels) the subhalo is moved from close to the cluster outwards and (second two panels) the mass fraction is increased ( $M_{\text{sub}}/M_{\text{main}}$ ).

The fitted lines show over what mass interval the mass-independent bias remains, until the chi-square of the line becomes greater than one standard deviation from the expected value. Fig. 8 shows that the fit breaks down at low radii ( $< 30$  arcsec). This value coincides with the size of the prior around the subhalo and shows that the sampler cannot de-merge the two haloes. Figs 8(c) and (d) show that the bias is independent of mass ratio, and even in the case where the two haloes are of equivalent size the bias remains



**Figure 8.** Weak lensing accuracy as a function of the radial position from the cluster and mass fraction ( $M_{\text{sub}}/M_{\text{main}}$ ). In each case the background galaxies have intrinsic ellipticities, redshift distribution and shape measurement bias. The first two panels (main and subhalo respectively) show an  $8 \times 10^{13} M_{\odot}$  cluster (with associated 10 times large parent halo), simulated at various distances from the main halo. The catastrophic failure at  $< 30''$  is due to the subhalo position overlapping with the parent halo. The second two panels show the positional estimates of a main halo of  $8 \times 10^{14}$  and subhalo with an increasing subhalo mass (decreasing ratio). It is shown that the bias is mass fraction independent and is robust to minor mergers as well as substructure infall. (a) Main Halo. (b) Sub Halo. (c) Main Halo. (d) Sub Halo.



**Figure 9.** Main (left) and Sub (right panel) position as function of cluster member contamination expressed as a fraction of the total background galaxy number.

negligible. This is promising as it shows the reconstruction should be reliable even in the case of a minor merger and not just substructure infall.

#### 4.2 Cluster member inclusion

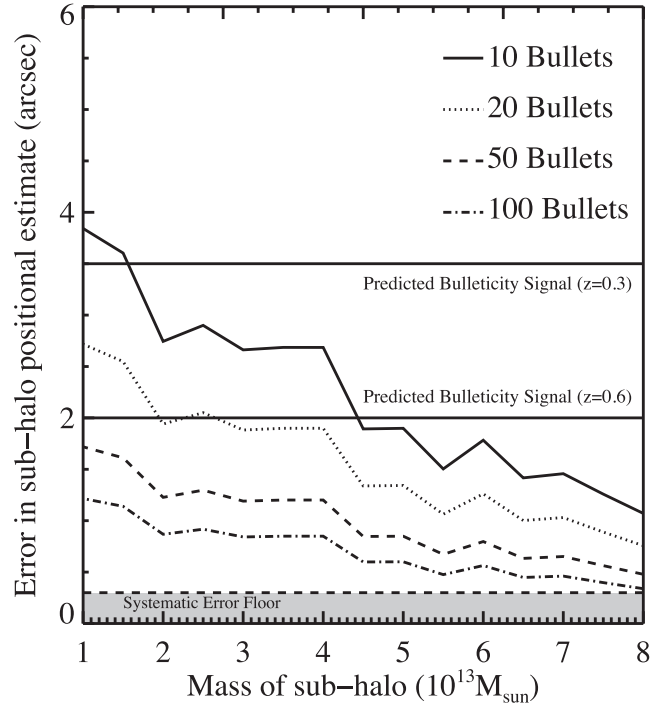
Fig. 9 shows the results from including member galaxies into the reconstruction. The fraction of background galaxies is calculated by summing the total number of member galaxies in the field of view and dividing by the number of background galaxies. It can be seen that the reconstruction is reliable up to  $\sim 30$  per cent of the background galaxies, at which the position of the subhalo and main halo becomes unreliable. Given a background density of  $80 \text{ arcmin}^{-2}$ , which is significantly less found by Hoekstra et al. (2011), we can conclude from these plots that we not worried about inclusion of these member galaxies.

### 5 PRECISION OF LENSTOOL

Throughout this investigation we have consistently found that by averaging many clusters together one can measure the position of DM haloes accurately to within  $\sim 0.3$  arcsec. Understanding how precisely we can measure these offsets informs us how many offsets will need to be measured in order to make a statistically significant detection.

The error bars derived from LENSTOOL give a rough estimate of the number of subhaloes that are required to robustly measure a significant offset between DM and gas. We decide to use the error bars from Fig. 5 to derive the precision. These were the error bars in the case of all signal contaminants. We found no evidence for additional uncertainty due to member galaxies at the expected level, and therefore have not factored these into the errors.

Dietrich et al. (2012) find that cosmic shear can offset the position of weak lensing peaks of order 5 kpc ( $0.7$  arcsec at a  $z = 0.6$ ). We expect such a contaminant to average out to zero but have a contribution to the overall error and precision. We therefore add this contaminant in quadrature to the error bars given in and calculate for a given size of clusters of a given mass with a given mass subhalo



**Figure 10.** The error in the mean position for various sample sizes. In order to detect an overall offset between baryonic and DM the error in the mean of the sample size needs to be less than the expected signal. For clusters at redshift  $z = 0.6$ , the bulleticity is  $\sim 2$  arcsec, and for a redshift  $z = 0.3$  this increases to  $3.5$  arcsec. A sample of  $\sim 50$  offsets should yield a significant detection of bulleticity, and LENSTOOL can measure these offsets with subdominant systematic bias.

for a cluster at redshift  $z = 0.6$ . Fig. 10 shows between  $\sim 20$  and  $50$  measured offsets are required in order to have statistically significant detection. In this scenario there is also shape measurement bias and a source galaxy redshift distribution.

Since typical clusters each have conservatively one in-falling group of galaxies containing  $\sim 10$  per cent of their mass, this suggests between  $20$  and  $50$  clusters are needed. This is feasible within

the current *HST* and *Chandra* archive. Furthermore, any strong lensing detection would tighten the constraints on the mass and concentration of the main halo can be more tightly constrained, and will lead to a better measurement of the offset. It is would be trivial to include a strong lensing model in to LENSTOOL, however, relative weighting of the constraints provided is an issue which will be addressed in future work.

Although basing predicted sample sizes on controlled environments such as those studied here, the results give us optimism to carry out the measurement in real data.

In the case of definite peaks and well-defined profiles we expect a sample of between 20 and 50 to be sufficient in order to measure a significant offset. We use this sample size as a confirmation that one should be able to make a detection using the current Hubble archive. We expect the true sample size to be larger than this and using hydrodynamical simulations, a more accurate sample size can be determined. Such tests are beyond the scope of this paper and will be carried out in conjunction with the data.

## 6 DISCUSSION

Through carefully controlled experiments, it was found that the likelihood surfaces for the reconstructed subhalo positions are symmetric around the true value in the regime of infinite signal-to-noise ratio. In the presence of trivial noise contaminants the estimated positions are also not biased, however, adding shape measurement bias seems to introduce a small bias of order  $\sim 0.3$  arcsec in the subhalo.

In order to better constrain our errors and any bias in position we averaged each mass scale over each single subhalo configuration, and dual halo configuration. It was found that the positional bias in all cases is independent of mass and configuration. For a single halo configuration the only cause of bias was due to imperfect shape measurement. The observed offset of  $0.27 \pm 0.14$  is well within our tolerated level. This bias was seen throughout the simulation using NFW profiles, including the dual subhalo configurations. In the case of an SIS profile, we found that the positional estimate no longer observed a bias, which was due to the peaky nature of the central core, however, the errors become unacceptably large below  $1.5 \times 10^{13} M_{\odot}$ .

In all of the simulations we do not find strong evidence for a positional bias of greater than 0.5 arcsec. Importantly, this performance is sufficient to enable a detection of the theoretically expected  $\sim 2.0$  arcsec ( $\sim 3.5$  arcsec) offset between substructure's DM and gas as it falls into massive clusters at a redshift of 0.6 (0.3) (MKN).

Initial work using the *HST* archive will aim to measure the average displacement between DM and both gas and stars, that could provide evidence for self-interacting DM. This displacement can then be calibrated with simulations of interacting DM to estimate its cross-section.

With a sample of clusters already available in the *HST* archive, averaging the measured offset of many pieces of substructure will provide sufficiently accurate DM centroiding to detect an offset. Future space missions (e.g. *Euclid*;<sup>2</sup> Laureijs et al. 2011) will increase the sample of available clusters by many orders of magnitude. However, fully exploiting such data would require improved mass reconstruction techniques

Furthermore, in the quasi-weak lensing regime considered here, flexion, the third derivative of the lensing potential, becomes im-

portant. As the gradient of the tidal field it is more sensitive to small-scale structure, similar to that investigated here (Bacon et al. 2006). The positional information in this higher order process could therefore provide significantly improved offset measurements. Unfortunately, flexion remains extremely difficult to measure and it fundamental properties such as the intrinsic flexion distribution and accuracy requirements are still yet to be determined (Viola, Melchior & Bartelmann 2012). Future algorithms could potentially exploit this extra information, however, this is currently not possible.

## 7 CONCLUSIONS

Measuring the separation of DM and baryonic gas in groups of infalling galaxies requires accurate astrometry. In the introduction we proposed five primary questions.

- (i) To what precision can we measure the offset in a single infalling substructure?
- (ii) Can we identify a point estimator whose simple combination over a sample will provide a measurement of subhalo position with minimal bias?
- (iii) What are the dominant sources of residual bias in this estimate?
- (iv) How large a sample of observed clusters are we likely to need to be able to detect an offset between DM and baryonic as predicted by MKN?
- (v) What further investigation might be needed to prove the utility of this technique for probing DM interaction cross-sections?

In reference to these, we find the following.

- (i) The public LENSTOOL software can measure the position of individual  $1.5 \times 10^{13} M_{\odot}$  peaks with  $\sim 0.3$  arcsec systematic bias, as long as they are at least  $\sim 30$  arcsec from the cluster centre. Any subhaloes detected above this threshold will be real and only biased to  $\sim 0.3$  arcsec.
- (ii) The maximum likelihood value of the two dimensional position likelihood surface is found to be the best point source estimator, being negligibly biased in the noise-free case compared to the mean value estimator.
- (iii) The dominant source of bias is caused by a preferred direction to the shape of galaxies introduced by a biased shape measurement algorithm.
- (iv) Since typical clusters each have on average one infalling groups of galaxies containing  $\sim 10$  per cent of their mass, between 20 and 50 clusters are needed to detect an offset between dark and baryonic matter.
- (v) The method will need to be tested on full hydrodynamical simulations (containing a more complex distribution of mass) in parallel with real data to show that the displacement obtained from data is reliable.

This work gives us confidence to pursue offsets as a technique in the measurement of the DM cross-section.

## ACKNOWLEDGEMENTS

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<sup>2</sup> <http://www.euclid-ec.org>

## REFERENCES

- Angloher G. et al., 2012, *Eur. Phys. J. C*, 72, A1971
- Bacon D. J., Goldberg D. M., Rowe B. T. P., Taylor A. N., 2006, *MNRAS*, 365, 414
- Baer H., Tata X., 2009, in Datta A., Mukhopadhyaya B., Raychaudhari A., eds, *Physics at the Large Hadron Collider*. Springer India, New Delhi, p. 179
- Bartelmann M., Schneider P., 2001, *Phys. Rep.*, 340, 291
- Battaner E., Florido E., 2000, *Fundamentals Cosmic Phys.*, 21, 1
- Bernabei R. et al., 2010, *European Phys. J. C*, 67, 39
- Bertone G., Hooper D., Silk J., 2005, *Phys. Rep.*, 405, 279
- Bradač M., Schneider P., Lombardi M., Erben T., 2005, *A&A*, 437, 39
- Bradač M. et al., 2006, *ApJ*, 652, 937
- Bradač M., Allen S. W., Treu T., Ebeling H., Massey R., Morris R. G., von der Linden A., Applegate D., 2008, *ApJ*, 687, 959
- Bridle S. et al., 2010, *MNRAS*, 405, 2044
- Burgos S. et al., 2009, *Astropart. Phys.*, 31, 261
- Cacciato M., Bartelmann M., Meneghetti M., Moscardini L., 2006, *A&A*, 458, 349
- Cholis I., Goodenough L., Hooper D., Simet M., Weiner N., 2009, *Phys. Rev. D*, 80, 123511
- Clowe D., Gonzalez A., Markevitch M., 2004, *ApJ*, 604, 596
- Clowe D., Bradač M., Gonzalez A. H., Markevitch M., Randall S. W., Jones C., Zaritsky D., 2006, *ApJ*, 648, L109
- Cohn J. D., 2012, *MNRAS*, 419, 1017
- Diego J. M., Tegmark M., Protopapas P., Sandvik H. B., 2007, *MNRAS*, 375, 958
- Dietrich J. P., Böhnert A., Lombardi M., Hilbert S., Hartlap J., 2012, *MNRAS*, 419, 3547
- Duffy A. R., Schaye J., Kay S. T., Dalla Vecchia C., Battye R. A., Booth C. M., 2010, *MNRAS*, 405, 2161
- Finoguenov A. et al., 2007, *ApJS*, 172, 182
- Heymans C. et al., 2006, *MNRAS*, 368, 1323
- Hoekstra H., 2001, *A&A*, 370, 743
- Hoekstra H., Jain B., 2008, *Annu. Rev. Nucl. Part. Syst.*, 58, 99
- Hoekstra H., Donahue M., Conselice C. J., McNamara B. R., Voit G. M., 2011, *ApJ*, 726, 48
- Hooper D., Goodenough L., 2011, *Phys. Lett. B*, 697, 412
- Jetzer P., Koch P., Piffaretti R., Puy D., Schindler S., 2001, *Proc. Symp., New Visions of the X-ray Universe in the XMM-Newton and Chandra era, Morphology of Galaxy Clusters and Sunyaev-Zel'dovich effect*, preprint (astro-ph/0201421)
- Jullo E., Kneib J.-P., Limousin M., Elíasdóttir Á., Marshall P. J., Verdugo T., 2007, *New J. Phys.*, 9, 447
- Kitching T. et al., 2012a, *MNRAS*, 423, 3163
- Kitching T. et al., 2012b, *MNRAS*, in press (arXiv:1204.4096)
- Lasky P. D., Fluke C. J., 2009, *MNRAS*, 396, 2257
- Laureijs R. et al., 2011, *Euclid Definition Study Report*, ESA/SRE(2011)12, preprint (arXiv:1110.3193)
- Leauthaud A. et al., 2007, *ApJS*, 172, 219
- Leauthaud A. et al., 2010, *ApJ*, 709, 97
- Macciò A. V., Dutton A. A., van den Bosch F. C., 2008, *MNRAS*, 391, 1940
- Mahdavi A., Hoekstra H., Babul A., Balam D. D., Capak P. L., 2007, *ApJ*, 668, 806
- Markevitch M., Gonzalez A. H., Clowe D., Vikhlinin A., Forman W., Jones C., Murray S., Tucker W., 2004, *ApJ*, 606, 819
- Massey R., Goldberg D. M., 2008, *ApJ*, 673, L111
- Massey R. et al., 2007a, *Nat*, 445, 286
- Massey R. et al., 2007b, *MNRAS*, 376, 13
- Massey R., Kitching T., Richard J., 2010, *Rep. Progress Phys.*, 73, 086901
- Massey R., Kitching T., Nagai D., 2011, *MNRAS*, 413, 1709 (MKN)
- Melchior P., Viola M., 2012, *MNRAS*, 424, 2757
- Merten J., Cacciato M., Meneghetti M., Mignone C., Bartelmann M., 2009, *A&A*, 490, 681
- Merten J. et al., 2011, *MNRAS*, 417, 333
- Mitsou V. A., 2011, *J. Phys.: Conf. Ser.*, 335, 012003
- Nagai D., Vikhlinin A., Kravtsov A., 2007a, *ApJ*, 655, 98
- Nagai D., Kravtsov A., Vikhlinin A., 2007b, *ApJ*, 668, 1
- Navarro J. F., Frenk C. S., White S. D. M., 1996, *ApJ*, 462, 563 (NFW)
- Peter A. H. G., Rocha M., Bullock J. S., Kaplinghat M., 2013, *MNRAS*, 430, 105
- Powell L. C., Kay S. T., Babul A., 2009, *MNRAS*, 400, 705
- Randall S. W., Jones C., Markevitch M., Blanton E. L., Nulsen P. E. J., Forman W. R., 2009, *ApJ*, 700, 1404
- Refregier A., 2003, *ARA&A*, 41, 645
- Rhodes J. et al., 2007, *ApJS*, 172, 203
- Rocha M., Peter A. H. G., Bullock J. S., Kaplinghat M., Garrison-Kimmel S., Onorbe J., Moustakas L. A., 2013, *MNRAS*, 430, 81
- Rubin V. C., Ford W. K. J., Thonnard N., 1980, *ApJ*, 238, 471
- Schmidt F., Leauthaud A., Massey R., Rhodes J., George M. R., Koekemoer A. M., Finoguenov A., Tanaka M., 2012, *ApJ*, 744, L22
- Scoville N. et al., 2007, *ApJS*, 172, 1
- Shan H. Y., Qin B., Zhao H. S., 2010, *MNRAS*, 408, 1277
- Spinelli P. F., Seitz S., Lerchster M., Brimiouille F., Finoguenov A., 2012, *MNRAS*, 420, 1384
- Taylor A. N., Kitching T. D., Bacon D. J., Heavens A. F., 2007, *MNRAS*, 374, 1377
- Viola M., Melchior P., Bartelmann M., 2012, *MNRAS*, 419, 2215
- Williams L. L. R., Saha P., 2011, *MNRAS*, 415, 448
- Wright C. O., Brainerd T. G., 2000, *ApJ*, 534, 34
- Yoshida N., Springel V., White S. D. M., Tormen G., 2000, *ApJ*, 544, L87
- Zitrin A., Bartelmann M., Umetsu K., Oguri M., Broadhurst T., 2012, *MNRAS*, 426, 2944
- Zwicky F., 1933, *Helvetica Phys. Acta*, 6, 110

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